



# Low-coordination transition metal sites on oxygen vacancy enriched strontium titanate-based perovskites enable highly selective photocatalytic CO<sub>2</sub> reduction to CH<sub>4</sub>

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## ABSTRACT

The photochemical conversion of CO<sub>2</sub> into high-value single hydrocarbon fuels such as CH<sub>4</sub> remains a challenge. Here we have prepared a series of perovskites V<sub>0</sub>-SrTiM<sub>0</sub><sub>3</sub> (M = Mn, Fe, Co) with rich oxygen vacancies. Among them, Mn-doped perovskite (V<sub>0</sub>-STMn<sub>0.2</sub>) exhibited 18.21 μmol g<sup>-1</sup> yield and nearly 100 % selectivity for CH<sub>4</sub>, as well as excellent reproducibility of about 15 h. Through a combination of advanced characterization and theoretical calculations, we found that the low-coordination transition metal site was able to stabilize the critical COOH\* and CO\* intermediates, thereby altering the reaction pathway to form CH<sub>4</sub> instead of CO. Most importantly, this work demonstrated that the relative center between the 2p orbitals of the oxygen atom of the CO<sub>2</sub> reduction intermediate and the 3d orbitals of the low-coordination metal site regulates the reaction mechanism of CO<sub>2</sub> reduction, offering the possibility of achieving efficient photochemical reduction of CO<sub>2</sub>.

## 1. Introduction

The direct use of solar energy to boost carbon recovery and convert CO<sub>2</sub> into hydrocarbon fuels is undoubtedly a promising strategy for mitigating the greenhouse effect and securing future energy supplies [1]. Unfortunately, photoreduction of the extremely stable CO<sub>2</sub> molecule is very difficult and elaborate owing to the high C=O dissociation energy of about 750 kJ mol<sup>-1</sup> [2]. Moreover, due to the complex electron and proton transfer processes involved, a wide variety of products such as CO, CH<sub>3</sub>OH and CH<sub>4</sub> can be produced in the CO<sub>2</sub> photoreduction process. Therefore, it remains a major challenge to design highly active photocatalytic systems with a highly pure gaseous product.

In recent years, researchers have tried to improve the light absorption and CO<sub>2</sub> adsorption/activation of photocatalysts by introducing foreign elements [3], co-catalysts loaded with noble metals [4], surface defects [5] and functional groups [6,7]. For perovskite oxides (ABO<sub>3</sub>), metal cation doping involves substitution of A or B sites, and the d-type orbitals of the doped metal cations will overlap and hybridize with those of the existing A/B site metals, and exhibit d-d transition sub-peak

absorption in the low photon energy region [8,9], which can dramatically alter the band structure and create new activity for photocatalytic CO<sub>2</sub> reduction [10]. SrTiO<sub>3</sub> is an archetypal perovskite oxide, however, it is not a good candidate for a photocatalyst since it is only active under ultraviolet radiation [11]. Many studies have investigated the introduction of various dopants such as Cr [12], Co [13], Rh [14], Al [15] and Bi [16] to extend the light absorption of SrTiO<sub>3</sub> into the visible range. In addition, the introduction of oxygen vacancies into perovskite oxides is also an effective approach to modify their photocatalytic performance [17]. Oxygen vacancies can reduce the coordination numbers of the neighboring metal atoms to provide more active sites for the intermediates [18,19], thus further optimizing energetics of the catalytic reaction [20,21]. In photocatalysis, low-valence electron-rich 3d transition metal ions with specific electronic structures (such as Ti<sup>3+</sup> (III) [22], Mn<sup>3+</sup> (III) [23], Ni<sup>+</sup> (I) [24], and Cu<sup>δ+</sup> ( $\delta < 2$ ) [25]) can act as catalytic centers or promoters to enhance the kinetics of a variety of thermodynamically unfavorable reactions [26]. Advanced structural and theoretical analyses attribute the excellent activity to Jahn-Teller distortion, large amounts of oxygen vacancies around the doped

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transition metal ions and the introduction of compressive strain into the host. This synergistic effect promotes molecular adsorption and activation, and rapid charge separation [27]. Although numerous theoretical and experimental investigations have been carried out on the surface defects of perovskite oxides [28], detailed studies on the optimization of the chemical absorption behavior of intermediates in the  $\text{CO}_2$  reduction process are still lacking.

Herein, based on defect chemistry theory, we constructed oxygen vacancy enriched perovskites  $\text{SrTiMO}_3$  ( $\text{M} = \text{Mn, Fe, Co}$ ) for  $\text{CO}_2$  photoreduction. Detailed characterization shows that the introduction of oxygen vacancies creates an abundance of low-coordination transition metal sites, which significantly improves the separation and transfer efficiency of carriers. In situ diffuse reflectance infrared Fourier transform spectroscopy (DRIFTS) and density functional theory (DFT) calculations confirm that the low-coordination transition metal sites can lower the energy barrier from  $\text{CO}^*$  to  $\text{CHO}^*$ . Meanwhile, the relative center between the 2p orbitals of the oxygen atom of the  $\text{CO}_2$  reduction intermediate and the 3d orbitals of the low coordination metal site regulates the adsorption stability of the low coordination metal site to the oxygen intermediate, ultimately conferring the selectivity for desired  $\text{CH}_4$  product.

## 2. Experimental

### 2.1. Chemicals and materials

Srontium nitrate ( $\text{Sr}(\text{NO}_3)_2$ , AR, Sinopharm Chemical Reagent), manganese chloride tetrahydrate ( $\text{MnCl}_2 \cdot 4 \text{H}_2\text{O}$ , 99.99 %, Macklin), iron nitrate nonahydrate ( $\text{Fe}(\text{NO}_3)_3 \cdot 9 \text{H}_2\text{O}$ , 99.99 %, Macklin), cobalt nitrate hexahydrate ( $\text{Co}(\text{NO}_3)_2 \cdot 6 \text{H}_2\text{O}$ , 99.99 %, Macklin), tetrabutyl titanate ( $\text{C}_{16}\text{H}_{36}\text{O}_4\text{Ti}$ , AR, Macklin), ethanolamine (EA, AR, Macklin), sodium hydroxide ( $\text{NaOH}$ , 95 %, Sinopharm Chemical Reagent), Sodium borohydride ( $\text{NaBH}_4$ , AR, Sinopharm Chemical Reagent). All the chemicals were used directly without any further purification.

### 2.2. Synthesis of $\text{SrTiO}_3$ and $\text{SrTi}_{1-x}\text{M}_x\text{O}_3$ perovskite oxides

The pure- $\text{SrTiO}_3$  and  $\text{SrTi}_{1-x}\text{M}_x\text{O}_3$  ( $\text{M} = \text{Mn, Fe, Co}$ ) were synthesized by a facile one-pot solvothermal method [29]. In a typical procedure, 8 mmol of tetrabutyl titanate and 2 mmol  $\text{MnCl}_2 \cdot 4 \text{H}_2\text{O}$  were dissolved in 20 mL of EA and stirred for 30 min. 50 mL of 3 M  $\text{NaOH}$  aqueous solution was added to the solution and stirred for a further 30 min to obtain the precipitate suspension. 10 mL of aqueous solution containing 10 mmol of  $\text{Sr}(\text{NO}_3)_2$  was added into the suspension and the solution stirred for 30 min. The solution was then transferred into a 100 mL Tefon-lined autoclave, sealed, and heated at 180 °C for 24 h, and allowed to cool to room temperature naturally. The final product was collected by centrifuging the mixture, washed with ethanol and deionized water many times until the organic residuals were completely removed, and then dried in vacuum at 80 °C for 10 h. The obtained  $\text{STMn}_{0.2}$  perovskite oxides were calcined at 350 °C for 2 h in static air and then cooled to room temperature for further characterization. By tuning the molar ratio of tetrabutyl titanate and transition metal ion precursor chemicals in the final suspension solution, the obtained  $\text{SrTi}_{1-x}\text{M}_x\text{O}_3$  samples were designed as  $\text{STM}_x$  ( $x = 0.1, 0.2, 0.3, 0.4$ ), and the pure- $\text{SrTiO}_3$  sample was termed STO.

### 2.3. Synthesis of oxygen vacancy enriched $\text{SrTiO}_3$ and $\text{SrTi}_{1-x}\text{M}_x\text{O}_3$ perovskite oxides

In a typical procedure, the obtained  $\text{SrTiO}_3$  ( $\text{SrTi}_{1-x}\text{M}_x\text{O}_3$ ) perovskite oxide was completely ground with  $\text{NaBH}_4$  (with a mass ratio of 1:1) and the mixture was then placed in a porcelain crucible. It was further calcined at 350 °C for 2 h in Ar (99.999 %) at a ramp rate of 2.5 °C  $\text{min}^{-1}$ , and allowed to cool to room temperature. Finally, the resulting powders were washed with ethanol and deionized water many times

until the organic residuals were completely removed, and then dried in vacuum at 80 °C for 10 h. The obtained powder was designed as  $\text{V}_0\text{-STO}$  ( $\text{V}_0\text{-STM}_x$ ) and collected for further characterization.

Additional experimental details including materials, characterization, parameters for DFT calculations, and the photocatalytic test process can be found in the [Supplementary Material](#).

## 3. Results and discussion

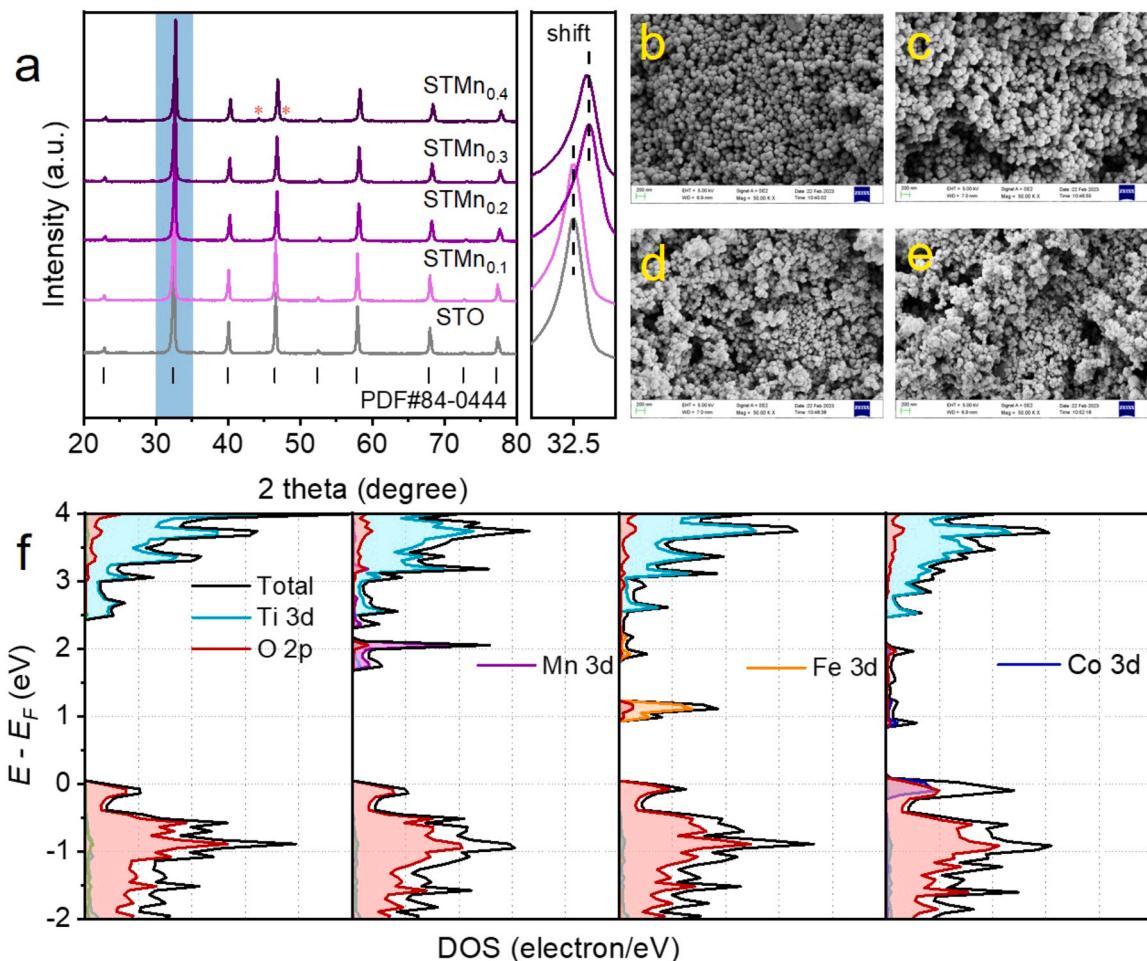
### 3.1. Structural and electronic states

As is displayed in [Fig. 1a](#) and [Fig. S1](#), the XRD patterns indicate that the transition metal ions  $\text{M}$  can be successfully doped into the perovskite STO lattice (PDF#84-0444) within the doping ratio range of 0.1–0.3. However, when the doping ratio is 0.4, the XRD pattern displays obvious impurity peaks at 40–50° (marked with \*), which can be indexed to  $\text{SrCO}_3$  (PDF#05-0418). Therefore, only the  $\text{STM}_x$  catalysts with  $\text{M}$ -doping ratio from 0.1 to 0.3 were investigated. For  $\text{STM}_{0.3}$ , the peaks were slightly shifted to higher angles compared to  $\text{STM}_{0.1}$ , indicating lattice shrinkage due to partial replacement of the larger  $\text{Ti}^{4+}$  ion (0.745 Å) by the smaller  $\text{Mn}^{3+}/\text{Mn}^{4+}$  (0.72/0.67 Å),  $\text{Fe}^{3+}$  (0.69 Å) or  $\text{Co}^{3+}$  (0.685 Å) ions [30]. Notably,  $\text{STMn}_{0.2}$  showed the same degree of peak shift as  $\text{STMn}_{0.3}$ . The Rietveld refinement method was used to fit the XRD results to obtain more detailed crystal structure information ([Fig. S2](#)). The refinement results reveal that perovskites STO,  $\text{STMn}_{0.1}$ ,  $\text{STMn}_{0.2}$  and  $\text{STMn}_{0.3}$  have the same structure of  $\text{Pm-3m}$ , and the unit cell volume decreases from 59.456 to 58.533 Å<sup>3</sup> after doping with Mn ions. The SEM results show that the average particle size of STO is around 140 nm ([Fig. 1b](#)). Following Mn ion doping, the average particle size of the  $\text{STM}_x$  gradually decreases ([Fig. 1c-e](#)).

The density of states (DOS) of STO before and after the substitution of the transition metal ion was further investigated. As shown in [Fig. 1f](#), the valence band (VB) of STO mainly consisted of O 2p orbitals and the conduction band (CB) primarily of Ti 3d states. The band gap of STO is about 2.4 eV, smaller than the experimental value because of the limitations of PBE method. After doping with M ions, the STO can be observed with a new energy level consisting mainly of M 3d and O 2p, resulting in a narrower band gap. For the  $\text{STM}_x$ , the new energy level as shallow donor energy levels near the bottom of the CB can absorb additional long wavelength visible light. The new energy levels of STFe and STCo as shallow acceptor energy levels near the top of the VB may also lead to an extension of the light absorption range. In conclusion, these impurity states can prolong the lifetime of the photogenerated carriers, thereby increasing the efficiency of the reaction.

After reduction with  $\text{NaBH}_4$ , XPS spectroscopy was carried out on the  $\text{V}_0\text{-STM}_x$  to determine its electronic states. [Fig. 2a](#) shows the O 1s spectra, which can be deconvoluted into peaks corresponding to the various oxygen species including lattice oxygen ( $\text{O}_{\text{lat}}$ , 529.3–530.5 eV), vacancy oxygen ( $\text{O}_v$ , 530.5–531.7 eV) and surface-adsorbed oxygen ( $\text{O}_{\text{ads}}$ , 531.8–532.8 eV), respectively [31,32]. Compared to  $\text{V}_0\text{-STMn}_{0.1}$ , the binding energy of O 1s of  $\text{V}_0\text{-STMn}_{0.2}$  shifts into the higher energy region by ~0.6 eV. The fitting results of O 1s spectra confirm that  $\text{V}_0\text{-STMn}_{0.2}$  (26.1 %) has a higher content of oxygen vacancies than  $\text{V}_0\text{-STMn}_{0.1}$  (17.5 %) and  $\text{V}_0\text{-STMn}_{0.3}$  (19.2 %). The Mn 2p spectrum shows two peaks at 653 and 643 eV ([Fig. 2b](#)), which belongs to Mn 2p<sub>3/2</sub> and Mn 2p<sub>1/2</sub>, respectively. The peaks at 646 eV correspond to the  $\text{Mn}^{2+}$  satellite feature [28]. For  $\text{V}_0\text{-STMn}_{0.2}$ , it exhibits a significantly greater proportion of the peaks at 646 eV than that of  $\text{V}_0\text{-STMn}_{0.1}$ / $\text{V}_0\text{-STMn}_{0.3}$  (16.1 % vs 15.3/15.6 %), evidencing more abundant  $\text{Mn}^{n+}$  species in  $\text{V}_0\text{-STMn}_{0.2}$ . Compared to  $\text{STMn}_{0.2}$ , there is an imperceptible shift in the binding energy of Sr 3d and Ti 2p in  $\text{V}_0\text{-STMn}_{0.2}$  ([Fig. S3](#)). In addition, the Ti 2p peak of  $\text{V}_0\text{-STMn}_{0.2}$  with Gaussian distribution determines the existence of  $\text{Ti}^{4+}$ , while the percentage of  $\text{Ti}^{3+}$  hardly changes.

The XRD pattern of  $\text{V}_0\text{-STMn}_{0.2}$  also possesses a similar cubic shape ([Fig. S4](#)), while its unit cell volume is enlarged to 59.249 Å<sup>3</sup> compared to



**Fig. 1.** (a) XRD patterns of perovskites  $\text{STMn}_x$  ( $x = 0.1\text{--}0.4$ ). SEM image of (b) STO, (c)  $\text{STMn}_{0.1}$ , (d)  $\text{STMn}_{0.2}$ , and (e)  $\text{STMn}_{0.3}$ . (f) The density of states of STO and STM (M = Mn, Fe, Co).

that of  $\text{STMn}_{0.2}$ . As displayed in the TEM images of  $\text{V}_0\text{-STMn}_{0.2}$  and  $\text{STMn}_{0.2}$  in Fig. 2c and Fig. S5, both samples have a well-crystallized structure with an average particle size of 93.93 nm for  $\text{V}_0\text{-STMn}_{0.2}$  and 88.36 nm for  $\text{STMn}_{0.2}$ . Moreover, the HRTEM image in Fig. 2d shows two interplanar spacing values of 0.278 nm and 0.389 nm with an angle of 45°, attributed to the (110) and (100) facets of  $\text{V}_0\text{-STMn}_{0.2}$ , respectively. Accordingly, it could be concluded that the exposed surface consists of (001) facets of  $\text{V}_0\text{-STMn}_{0.2}$ . Meanwhile, it is worth remarking that the significant stacking defect can be observed in Fig. 2d (as indicated by the white dash), which is caused by the narrowing of the lattice spacing after Mn ion substitution. The corresponding elemental mapping in Fig. 2e-i confirms the homogeneous distribution of Sr, Ti, Mn and O elements, and the elemental ratio is approximately Sr:Ti:Mn = 1:0.79:0.17 based on energy dispersive spectroscopy (Fig. S6).

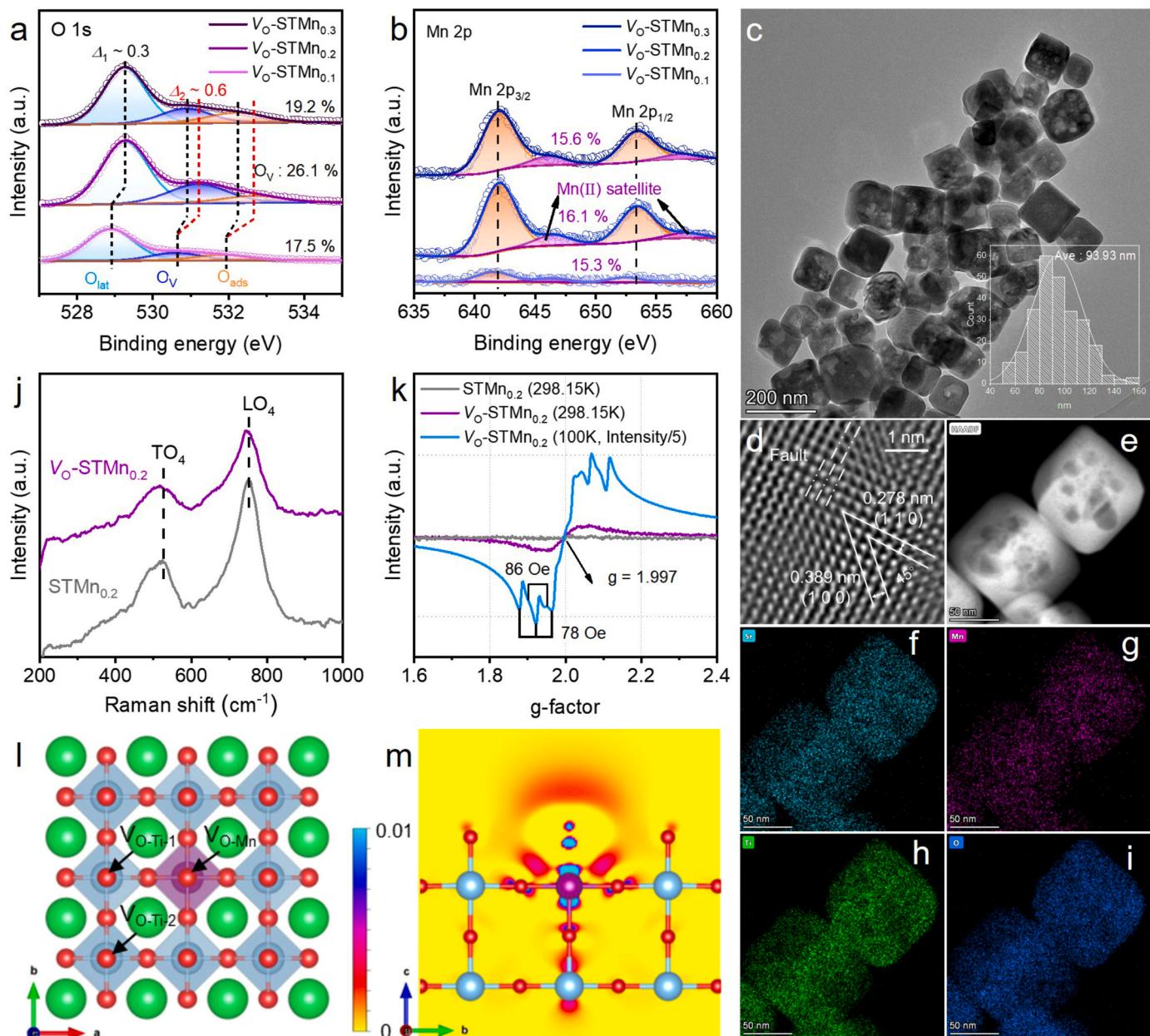
The structural information and oxidation state of the catalysts are revealed by Raman spectroscopy. As shown in Fig. 2j,  $\text{STMn}_{0.2}$  displays the typical  $\text{TO}_4$  and  $\text{LO}_4$  peaks at  $525\text{ cm}^{-1}$  and  $750\text{ cm}^{-1}$ , respectively [33]. In contrast,  $\text{V}_0\text{-STMn}_{0.2}$  shows a blueshift of the  $\text{TO}_4$  and  $\text{LO}_4$  peaks, indicating mass fluctuation at the oxygen position, which may be caused by the presence of oxygen vacancies. To provide more convincing information, electron spin resonance (ESR) spectroscopy was used to characterize the local environment of the present samples (Fig. 2k). The  $\text{STMn}_{0.2}$  displays a symmetrical ESR signal at  $g = 1.997$ , which is usually associated with the characteristics of  $\text{Mn}^{4+}$  ions. [34] After reduction, the  $\text{V}_0\text{-STMn}_{0.2}$  signal becomes stronger, indicating that more Mn ions are exposed due to the increased surface oxygen vacancies [35,36]. The ESR spectrum of  $\text{V}_0\text{-STMn}_{0.2}$  at 100 K was further

measured. As seen in Fig. 2j,  $\text{V}_0\text{-STMn}_{0.2}$  has a single sextet of lines with a hyperfine splitting of 78 Oe, which is indicative of  $\text{Mn}^{4+}$  ions doped at the Ti site in STO. Notably, in addition to the  $\text{Mn}^{4+}$  spectrum, a spectrum associated with  $\text{Mn}^{2+}$  appears with hyperfine splitting to 86 Oe. This obviously indicates that electrons supplied by oxygen vacancies are effectively trapped by the  $\text{Mn}^{4+}$  sites, converting  $\text{Mn}^{4+}$  to  $\text{Mn}^{2+}$  and maintaining the insulation of the sample [37].

It is reasonable to conclude that oxygen vacancies in perovskite  $\text{STMn}_{0.2}$  tend to be created in the  $\text{MnO}_6$  octahedra rather than the  $\text{TiO}_6$  octahedra. DFT computations were performed to further determine the oxygen vacancy formation energies ( $E_f$ ) of  $\text{STMn}_x$  on the  $\text{MnO}_6$  octahedron and  $\text{TiO}_6$  octahedron. The  $E_f$  is evaluated according to the following Eq. (1):

$$E_f = E_*^{V_0} + 1/2E_{O_2} - E_* \quad (1)$$

in which  $E_*^{V_0}$  is the total energy of the substrate structure with oxygen vacancy,  $E_*$  denotes the total energy of the substrate structure without oxygen vacancy, and  $E_{O_2}$  is the energy obtained from  $\text{O}_2$  molecule. As illustrated in Fig. 2l, the  $E_f$  values are 3.66 eV ( $\text{V}_0\text{-Mn}_0$ ) for the  $\text{MnO}_6$  octahedron and 4.45 eV ( $\text{V}_0\text{-Ti}_1$ ) or 4.37 eV ( $\text{V}_0\text{-Ti}_2$ ) for the  $\text{TiO}_6$  octahedron, suggesting that oxygen vacancies are more favorably generated in the  $\text{MnO}_6$  octahedron. Bader charge and differential charge density analysis was performed to study the charge redistribution and trapping (Fig. 2m). The creation of oxygen vacancies redistributes the charges on the surface as well as the layer beneath the surface of  $\text{STMn}_x$ , while the Mn atoms in the center of the  $\text{MnO}_6$  octahedron are allocated more electrons ( $+1.61\text{e} \rightarrow +1.16\text{e}$ ) than the Sr ( $+1.57\text{e} \rightarrow +1.52\text{e}$ ) atoms on

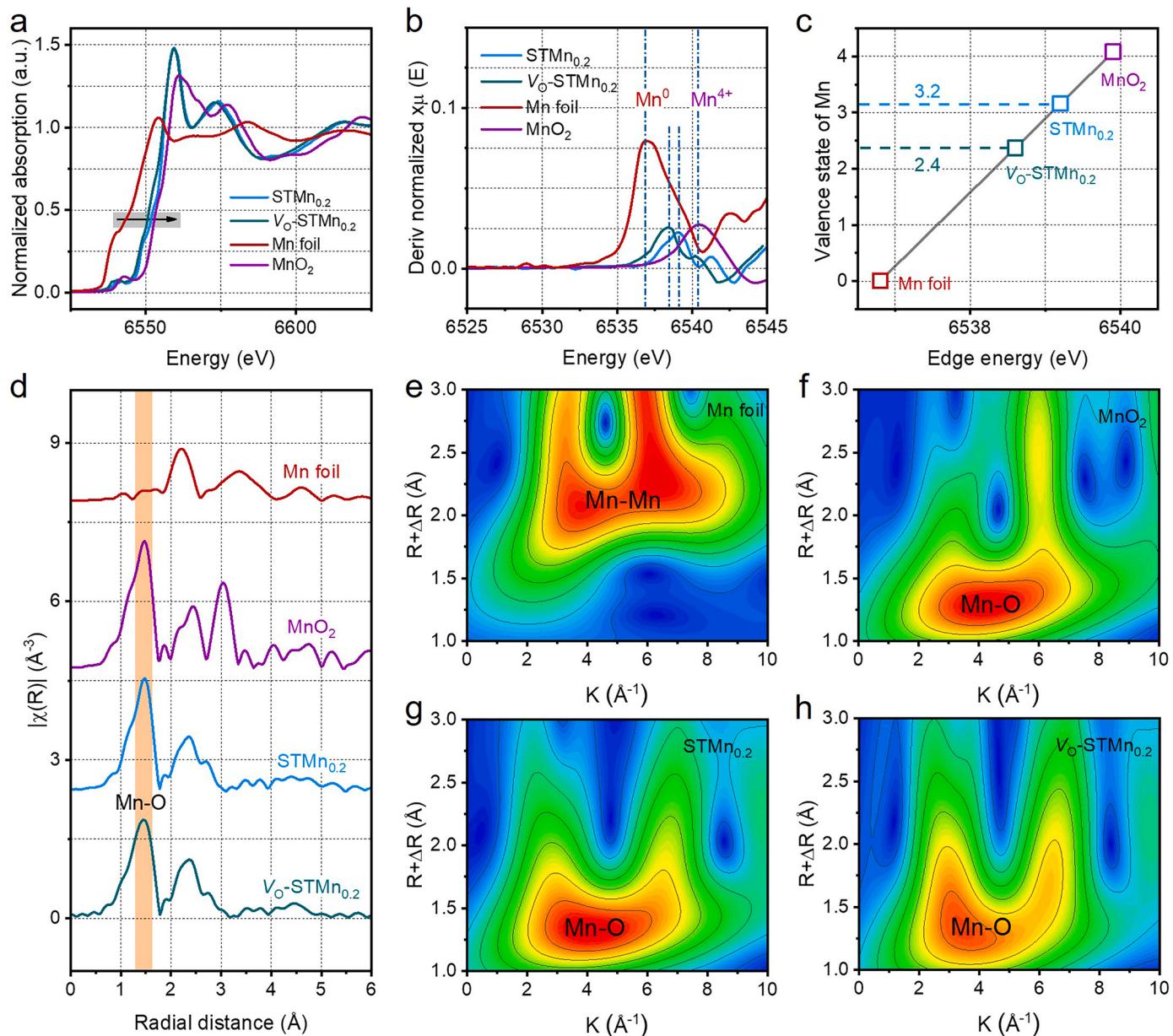


**Fig. 2.** The XPS spectra of (a) O 1s and (b) Mn 2p for  $V_O$ -STMn<sub>x</sub>. (c) TEM image of  $V_O$ -STMn<sub>0.2</sub>, the inset shows the particle size distribution with an average size of 93.93 nm. (d) HRTEM image of  $V_O$ -STMn<sub>0.2</sub>. (e-i) EDS elemental mappings of Sr (cyan), Mn (purple), Ti (green) and O (blue) with scale bars of 50 nm. (j) Raman shifts. (k) ESR spectra, in which the signal at  $g = 1.997$  corresponds to the Mn<sup>4+</sup> ions. (l) View of the positions of oxygen vacancy in STMn<sub>x</sub>. (m) 2D charge density differences. The charge density of blue represents electron accumulation, and yellow denotes electron depletion ( $e \text{ bohr}^{-3}$ ). The atomic color codes: green (Sr), cyan (Ti), purple (Mn), and red (O) balls, respectively.

the surface. Notably, these charge-enriched Mn sites may serve as highly active sites for strongly trapping C or O atoms into stable intermediate configurations, offering the possibility of selectively tailoring product.

XAS spectroscopy measurements at the Mn K-edges were carried out to assess the chemical environment of  $V_O$ -STMn<sub>0.2</sub> and STMn<sub>0.2</sub>. As presented in Figs. 3a and 3b, the normalized X-ray absorption near edge structure (XANES) spectra on the  $V_O$ -STMn<sub>0.2</sub> and its corresponding first derivative are obviously shifted to a lower energy direction compared with that of STMn<sub>0.2</sub>, manifesting the lower average Mn valence state. Based on the linear relationship between the peak positions of the first derivatives of the Mn XANES spectra and the oxidation states, we calculated the valence states of the Mn atoms (Fig. 3c) [38]. The  $V_O$ -STMn<sub>0.2</sub> shows the lower Mn valence states (2.4 +) than the STMn<sub>0.2</sub> (3.2 +), which is in accordance well with the XPS data. The extended

X-ray absorption fine structure (EXAFS) spectra provided further structural information (Fig. 3d). The peak between 1.0 and 2.0 Å in the R-space represents the Mn-O shell. However, the intensity of the peak for  $V_O$ -STMn<sub>0.2</sub> is weaker than that for the STMn<sub>0.2</sub>, confirming the reduced Mn-O coordination number (Fig. S7). Furthermore, least squares curve parameter fitting was performed to analyze the quantitative coordination numbers of Mn confined in the  $V_O$ -STMn<sub>0.2</sub> and STMn<sub>0.2</sub>. The results of the EXAFS data fitting analysis are summarized in Fig. S8 and Table S1. The coordination number for Mn-O decreased from 5.65 to 4.41 after reduction with NaBH<sub>4</sub>, which is attributed to the loss of lattice oxygen atoms. The corresponding EXAFS wavelet transform (WT) analysis results are shown in Fig. 3e-h. A smaller intensity maximum at about 4 Å<sup>-1</sup> is observed, which correlates with the Mn-O pathway in  $V_O$ -STMn<sub>0.2</sub>, illustrating the low Mn-O coordination configuration. It is



**Fig. 3.** (a) Mn K-edge XANES spectra. (b) The first derivation of Mn K-edge XANES. (c) Mn oxidation state. (d) FT  $k^2$ -weighted EXAFS spectra of  $\text{STMn}_{0.2}$ ,  $\text{V}_0\text{-STMn}_{0.2}$ , and standard samples. (e-h) WT analysis at the Mn K edge of different samples.

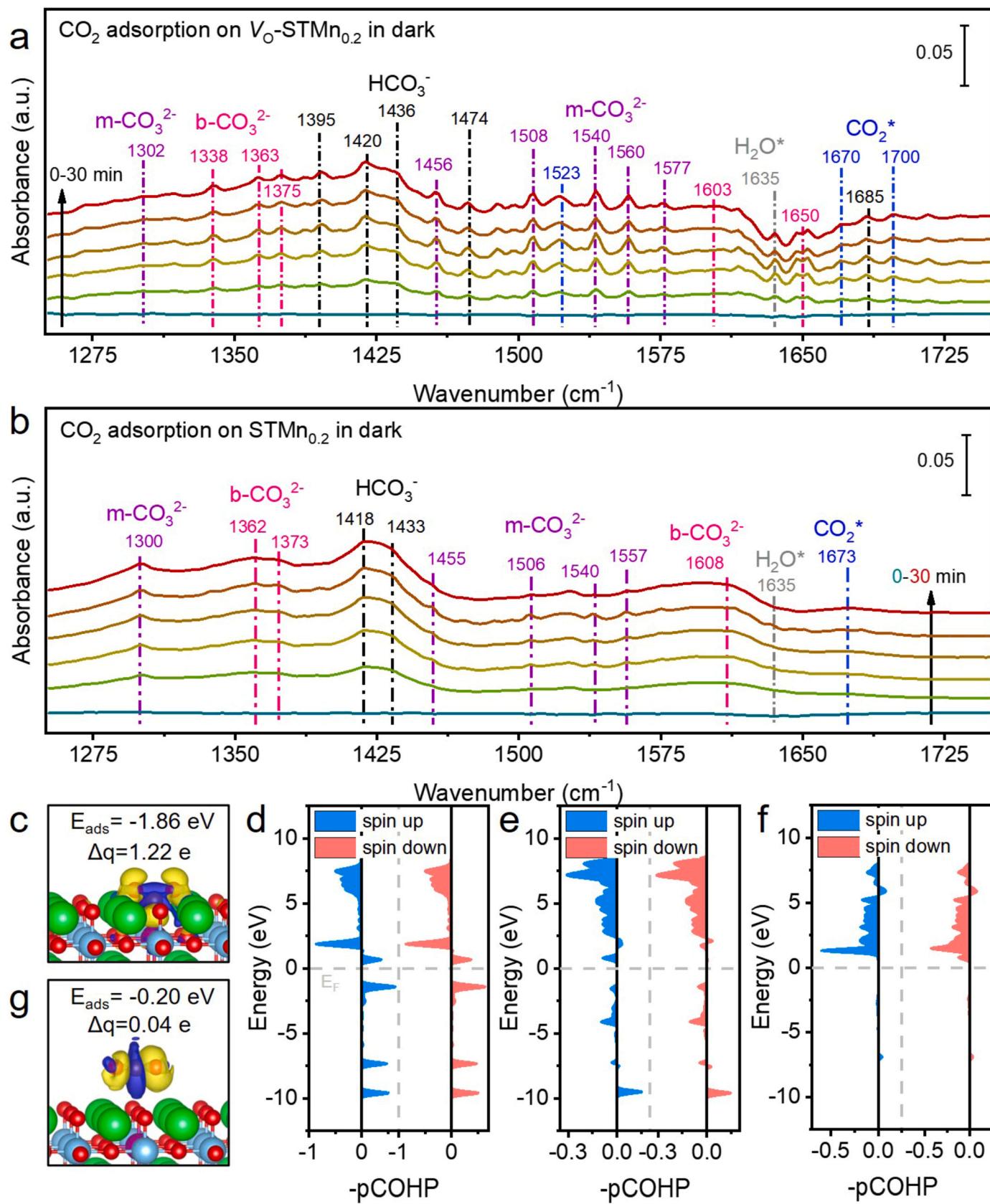
concluded that the lower coordination environment effectively regulates the electron state, and the altered  $e_g$  electron occupancy of Mn site can provide high active sites for catalytic reactions [19].

### 3.2. Photocatalytic performance for $\text{CO}_2$ reduction

As a precondition for photocatalytic  $\text{CO}_2$  reduction, the adsorption and activation of  $\text{CO}_2$  and  $\text{H}_2\text{O}$  on  $\text{STMn}_{0.2}$  and  $\text{V}_0\text{-STMn}_{0.2}$  surfaces in dark condition was studied using *in situ* DRIFTS. As depicted in Fig. 4a, exposure of  $\text{V}_0\text{-STMn}_{0.2}$  to the  $\text{H}_2\text{O}/\text{CO}_2$  atmosphere in the dark can induce the formation of highly adsorbed  $\text{H}_2\text{O}^*$  ( $1635\text{ cm}^{-1}$ ) and carbonate or bicarbonate species (monodentate ( $\text{m-}\text{CO}_3^{2-}$ ) at  $1302$ ,  $1458$ ,  $1508$ ,  $1540$ ,  $1560$  and  $1577\text{ cm}^{-1}$ , bidentate ( $\text{b-}\text{CO}_3^{2-}$ ) at  $1338$ ,  $1363$ ,  $1375$ ,  $1603$  and  $1650\text{ cm}^{-1}$ , and bicarbonate ( $\text{HCO}_3^-$ ) at  $1395$ ,  $1420$ ,  $1436$ ,  $1474$  and  $1685\text{ cm}^{-1}$ ) [39,40]. In particular, the key precursors of  $\text{CO}_2^*$  species are observed at  $1670$  and  $1700\text{ cm}^{-1}$ , indicating that the single-electron reduction of  $\text{CO}_2$  to  $\text{CO}_2^*$  can readily occur on the defective surface of  $\text{V}_0\text{-STMn}_{0.2}$ . Albeit of similar carbon active species

observed for  $\text{STMn}_{0.2}$ , the detected absorption peaks of these intermediates are significantly lower than those of  $\text{V}_0\text{-STMn}_{0.2}$  in the initial  $0\text{--}30\text{ min}$  of  $\text{CO}_2$  adsorption (Fig. 4b). For example,  $\text{STMn}_{0.2}$  produces dominant  $\text{m-}\text{CO}_3^{2-}$  instead of  $\text{b-}\text{CO}_3^{2-}$ , while the latter was the more robust site required for  $\text{CO}_2$  adsorption [18]. This is further supported by temperature programmed desorption (TPD) measurements of  $\text{CO}_2$ . Fig. S9 shows the  $\text{CO}_2$ -TPD curves of  $\text{STMn}_{0.2}$  and  $\text{V}_0\text{-STMn}_{0.2}$ , with the amount of  $\text{CO}_2$  adsorbed derived from the combined area of the curves. Clearly,  $\text{V}_0\text{-STMn}_{0.2}$  has a higher adsorption capacity compared to  $\text{STMn}_{0.2}$ , indicating that  $\text{V}_0\text{-STMn}_{0.2}$  has more  $\text{CO}_2$  adsorption and activation sites. As a result, the abundance of rationally created surface oxygen vacancies provides more adsorption sites and enhances the reactivity of  $\text{V}_0\text{-STMn}_{0.2}$ .

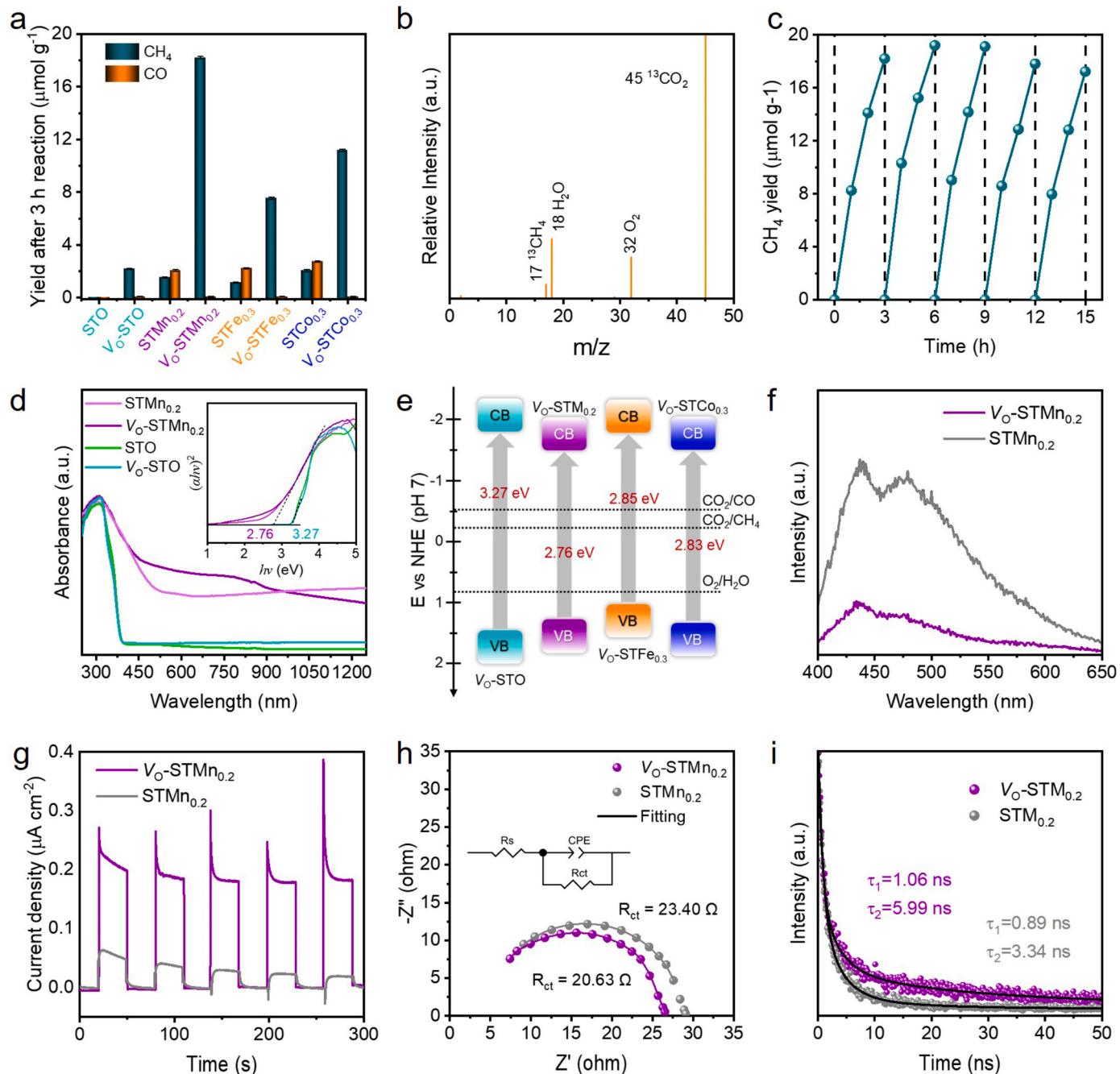
To better understand the adsorption and activation of  $\text{CO}_2$  on the surface with oxygen vacancies, DFT calculations were conducted to compare the  $\text{CO}_2$  adsorption over the oxygen vacancy and the perfect  $\text{STMn}_x$  surface. The adsorption energy ( $E_{\text{ads}} = \Delta G_{\text{CO}_2^*}$ ) of  $\text{CO}_2$  adsorption over the  $\text{V}_0\text{-STMn}_x$  ( $-1.86\text{ eV}$ ) is higher than that of  $\text{CO}_2$  adsorption on



**Fig. 4.** In situ DRIFTS test for  $\text{CO}_2$  and  $\text{H}_2\text{O}$  interaction with (a)  $\text{V}_\text{O}-\text{STMn}_{0.2}$  and (b)  $\text{STMn}_{0.2}$  in dark. Charge difference distributions for (c)  $\text{V}_\text{O}-\text{STMn}_\text{x}$  and (g)  $\text{STMn}_\text{x}$  following  $\text{CO}_2$  adsorption (charge accumulation is in yellow and depletion in blue, positive values for  $\Delta q$  indicate electron accumulation on  $\text{CO}_2$ , and  $E_{\text{ads}}$  is  $\text{CO}_2$  adsorption energy on surface). Isosurfaces are  $0.002 \text{ e bohr}^{-3}$ . Strontium, titanium, manganese, and oxygen atoms are denoted as green, cyan, purple, and red balls, respectively. (d) pCOHP between carbon atom in  $\text{CO}_2$  and Mn active site on  $\text{V}_\text{O}-\text{STMn}_\text{x}$ . pCOHP between oxygen atom in  $\text{CO}_2$  and surface Sr atom on (e)  $\text{V}_\text{O}-\text{STMn}_\text{x}$  and (f)  $\text{STMn}_\text{x}$ .

$\text{STMn}_x$  ( $-0.20$  eV). The bond situation of  $\text{CO}_2$  adsorbed in two different surfaces was then further compared by the charge difference and Bader charge analysis. As seen in Fig. 4c, the C atom is bound with the unsaturated Mn sites in  $\text{V}_0\text{-STMn}_x$ , causing the  $\text{O}=\text{C}=\text{O}$  bond to bend ( $122^\circ$ ) and elongate ( $1.28$  Å). The strong electron scale is distributed within the valence bonds, indicating the formation of a firm chemical bond between C atom and Mn site. The feedback  $\pi$  electrons from the Mn d orbital were transferred to the vacant orbital of C, resulting in an increase in the charge density of the C atom. The correlation results were confirmed by the Bader charge, where electrons ( $1.22$  e) were

transferred from the surface to the  $\text{CO}_2$  molecule. The above results were further supported by the calculation of the projected Crystal Orbital Hamilton Population (pCOHP) adsorbed by  $\text{CO}_2$  on  $\text{V}_0\text{-STMn}_x$  and  $\text{STMn}_x$ . In Fig. 4d, it is seen that most of the bonding orbitals of the C atom of  $\text{CO}_2$  and the Mn site are in the valence band below the Fermi level ( $E_F$ ), while the anti-bonding orbitals are mainly distributed in the conduction band above the  $E_F$ . In contrast, there are many anti-bonding orbitals below the  $E_F$  for O atom of  $\text{CO}_2$  and the surface Sr atom (Fig. 4e). Integral crystal orbital Hamiltonian (ICOHP) results show that the interaction between C atom and Mn active site is stronger than that



**Fig. 5.** (a) Production rates of  $\text{CH}_4$  (green columns) and  $\text{CO}$  (yellow columns) of  $\text{CO}_2$  photoreduction over  $\text{STO}$ ,  $\text{V}_0\text{-STO}$ ,  $\text{STMn}_x$  and  $\text{V}_0\text{-STMn}_x$  under Xenon lamp illumination. (b) Mass spectra of  $^{13}\text{CH}_4$  ( $\text{m/z} = 17$ ) and  $^{13}\text{CO}$  ( $\text{m/z} = 29$ ) produced in the photocatalytic reduction of  $^{13}\text{CO}_2$  over  $\text{V}_0\text{-STMn}_{0.2}$ . (c) Cycling measurements of  $\text{CO}_2$  photoreduction to  $\text{CH}_4$  for the  $\text{V}_0\text{-STMn}_{0.2}$  perovskite. (d) UV-vis diffuse reflectance spectra and inset for bandgap. (e) Schematic diagram of electronic band structures; the arrows represent the electron leap from the valence band (VB) to the conduction band (CB). (f) PL spectra for  $\text{V}_0\text{-STMn}_{0.2}$  and  $\text{STMn}_{0.2}$ . (g) Transient photocurrent density for  $\text{V}_0\text{-STMn}_{0.2}$  and  $\text{STMn}_{0.2}$  in 0.1 M  $\text{Na}_2\text{SO}_4$  aqueous solution. (h) EIS spectra for  $\text{V}_0\text{-STMn}_{0.2}$  and  $\text{STMn}_{0.2}$ . Inset is the fitted circuit. (i) TSPL spectra for  $\text{V}_0\text{-STMn}_{0.2}$  and  $\text{STMn}_{0.2}$ .

between O atom and surface Sr site (-ICOHP=2.34 for C-Mn vs -ICOHP=1.43 for O-Sr) [41]. Additionally, the weak interaction between CO<sub>2</sub> and the STMn<sub>x</sub> surface was also confirmed (-ICOHP=1.62 of O-Sr) and few electrons (0.04 e) transferred from the surface to the CO<sub>2</sub> molecule (Figs. 4f and 4g). In conclusion, these experimental and computational results demonstrate the stronger adsorption of CO<sub>2</sub> on V<sub>O</sub>-STMn<sub>0.2</sub> than on STMn<sub>0.2</sub>.

The photocatalytic CO<sub>2</sub> reduction performance of the samples was evaluated under xenon lamp illumination in a gas-solid reaction system (Scheme S1 and Fig. 5a). STO showed negligible CO<sub>2</sub> reduction activity, while V<sub>O</sub>-STO exhibited a yield of 2.17 μmol g<sup>-1</sup> CH<sub>4</sub>. In sharp contrast, the CO<sub>2</sub> conversion activity of V<sub>O</sub>-STM<sub>x</sub> was significantly enhanced after reduction with NaBH<sub>4</sub>. The STMn<sub>0.2</sub> displayed limited CO<sub>2</sub> photoreduction with 2.07 μmol g<sup>-1</sup> CO and 1.52 μmol g<sup>-1</sup> CH<sub>4</sub>. In comparison, V<sub>O</sub>-STMn<sub>0.2</sub> yielded 18.21 μmol g<sup>-1</sup> CH<sub>4</sub> with almost 100 % selectivity for CH<sub>4</sub> production, which demonstrated obvious superiority over the other Mn-doped samples (V<sub>O</sub>-STMn<sub>0.1</sub> 5.20 μmol g<sup>-1</sup> and V<sub>O</sub>-STMn<sub>0.3</sub> 9.36 μmol g<sup>-1</sup>, Fig. S10a). Meanwhile, the performance of V<sub>O</sub>-STFe<sub>x</sub> and V<sub>O</sub>-STCo<sub>x</sub> were boosted after active surface regulation, whereby the highest CH<sub>4</sub> yields of 7.55 μmol g<sup>-1</sup> and 11.19 μmol g<sup>-1</sup> were delivered by V<sub>O</sub>-STFe<sub>0.3</sub> and V<sub>O</sub>-STCo<sub>0.3</sub> (Figs. S11 and S12), respectively. Compared with other single component Ti-based photocatalysts, V<sub>O</sub>-STM<sub>x</sub> perovskites exhibit excellent CO<sub>2</sub> photoreduction performance under the absence of any co-catalysts and sacrificial reagents (Table S2). Moreover, O<sub>2</sub> evolution was detected in the photocatalytic reaction (Fig. S13). For V<sub>O</sub>-STMn<sub>0.2</sub>, the O<sub>2</sub> yield was 35.14 μmol g<sup>-1</sup>, approximately twice that of CH<sub>4</sub> yield (18.21 μmol g<sup>-1</sup>). The CH<sub>4</sub>/O<sub>2</sub> ratio was close to the theoretical value of 1:2, indicating that photogenerated holes were involved in the oxidation reaction of water to facilitate O<sub>2</sub> evolution.

For further verification, a series of control experiments were carried out, the results of which are shown in Fig. S10b. Negligible photocatalytic products were detected, whether the experiments were carried out in argon, in the dark, without H<sub>2</sub>O or without samples. Furthermore, the <sup>13</sup>CO<sub>2</sub> labelling experiment revealed that <sup>13</sup>CH<sub>4</sub> (*m/z* = 17) was the only CO<sub>2</sub> reduction product (Fig. 4b). Almost no H<sub>2</sub> (*m/z* = 2), <sup>13</sup>CO (*m/z* = 29) and <sup>13</sup>CH<sub>3</sub>OH (*m/z* = 33) were detected. The results above confirm that the evolution of CH<sub>4</sub> is exclusively originated from the photoreduction of CO<sub>2</sub>. Furthermore, surface oxygen vacancies are effective strategies for modulating electronic structure and adjusting surface properties to enhance photocatalytic activity. However, oxygen vacancies are prone to being filled by oxygen molecules or blocked by intermediate species, resulting in a gradual decrease in activity until catalyst deactivation occurs. Therefore, ensuring the stability of defects is a necessary condition for highly efficient and stable photocatalysts [42]. After 15 h of cyclic reactions, the CO<sub>2</sub> reduction performance of V<sub>O</sub>-STMn<sub>0.2</sub> only showed a slight decline compared to the initial state (17.23 vs 18.21 μmol g<sup>-1</sup>), as seen in Fig. 5c. Corresponding TEM, XRD, Raman, and XPS tests were conducted after 15 h of photocatalysis and are presented in Fig. S14. Considering these results comprehensively, our conclusion is that the loss of surface oxygen induces local lattice mismatch and forms a dense dislocation layer, which is clearly observed in Fig. 2d-e and Fig. S14a. As a result of the regulated coordination state and electronic structure of surface atoms, the dislocation layers can enhance the tolerability of oxygen-containing environments within a certain temperature range, thereby improving the cycle stability of the material [43,44].

### 3.3. Charge transfer mechanism

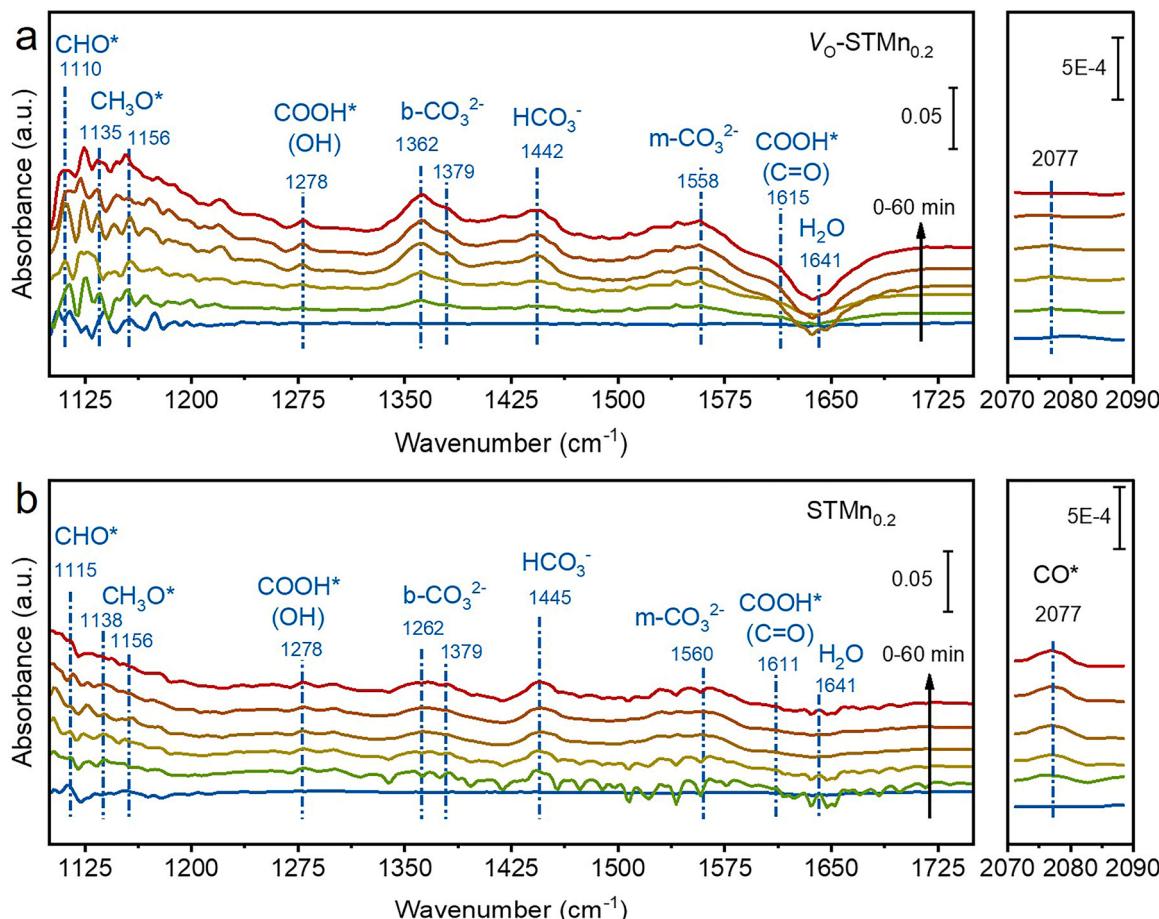
To determine the actual effects of transition metal substitution and oxygen vacancy formation, further optical measurements were carried out. The UV-Vis-NIR spectra (Fig. 5d) shows that STO and V<sub>O</sub>-STO exhibit the absorption onset at 390 nm, which agrees with the band gap edge absorption of SrTiO<sub>3</sub> (3.25 eV), whereas the additional absorption band beyond 600 nm is attributed to the increasing oxygen vacancy

concentration on V<sub>O</sub>-STO surface. For STMn<sub>0.2</sub>, due to the doped Mn changing the band gap structure, the band edge position at 480 nm indicates a narrower band gap. The optical response of V<sub>O</sub>-STMn<sub>0.2</sub> was enhanced in the spectrum range of 450–900 nm, which is ascribed to the more abundant oxygen vacancies. Notably, compared to Fig. 5a, the enhancement by visible light absorption cannot be neglected (Fig. S10b,  $\geq 400$  nm). A similar enhanced optical response in the visible light region was also observed in the V<sub>O</sub>-STFe<sub>0.3</sub> and V<sub>O</sub>-STCo<sub>0.3</sub> materials (Fig. S15a). Accordingly, the band gap of V<sub>O</sub>-STO, V<sub>O</sub>-STMn<sub>0.2</sub>, V<sub>O</sub>-STFe<sub>0.3</sub>, and V<sub>O</sub>-STCo<sub>0.3</sub> are calculated to be 3.27, 2.76, 2.85, and 2.83 eV, respectively. The positions of the VB maximum for V<sub>O</sub>-STO, V<sub>O</sub>-STMn<sub>0.2</sub>, V<sub>O</sub>-STFe<sub>0.3</sub>, and V<sub>O</sub>-STCo<sub>0.3</sub> were detected by XPS valence spectra (Fig. S15b). The corresponding standard hydrogen electrode (E<sub>VB</sub>, NHE) can then be calculated via the following formula: E<sub>VB</sub>, NHE =  $\varphi$  + E<sub>VB</sub>, XPS - 4.44, where  $\varphi$  is the work function of the instrument (4.2 eV). Thus, based on those results the electronic band structures versus NHE at pH = 7 can be elucidated and displayed in Fig. 5e. The results show that all the prepared perovskite oxides have suitable band gaps and band edge positions to realize CO<sub>2</sub> reduction and O<sub>2</sub> evolution simultaneously. Furthermore, the apparent quantum efficiency (AQE) for CH<sub>4</sub> evolution was evaluated under the illumination of monochromatic light. As exhibited in Fig. S15c, the AQE response of V<sub>O</sub>-STMn<sub>0.2</sub> matches well to the recorded UV-Vis spectrum, achieving AQE of 0.35 % at 380 nm.

Steady-state photoluminescence (PL) spectroscopy, transient photoluminescent (TSPL) spectroscopy, transient photocurrent (TPC) density and electrochemical impedance spectroscopy (EIS) were executed to probe the efficiency of charge-carrier separation and transportation. As shown by steady-state PL spectra, TPC responses and EIS spectra (Fig. 5f-h), the V<sub>O</sub>-STMn<sub>0.2</sub> has lower PL peak intensity, higher TPC density as well as smaller Nyquist semicircle radius in comparison to the STMn<sub>0.2</sub>. After fitting of the TSPL curves for Fig. 5i, prolonged carrier lifetimes ( $\tau_1$  = 1.06 ns;  $\tau_2$  = 5.99 ns) were observed for V<sub>O</sub>-STMn<sub>0.2</sub> compared to those for STMn<sub>0.2</sub> ( $\tau_1$  = 0.89 ns;  $\tau_2$  = 3.34 ns). In addition, the electronic structures of V<sub>O</sub>-STFe<sub>0.3</sub> and V<sub>O</sub>-STCo<sub>0.3</sub> were further analyzed. As shown in Fig. S16, V<sub>O</sub>-STMn<sub>0.2</sub> has the highest oxygen vacancy content, highest photocurrent density and lowest charge transfer resistance, even when compared to V<sub>O</sub>-STFe<sub>0.3</sub> and V<sub>O</sub>-STCo<sub>0.3</sub>. Moreover, the carrier lifetimes ( $\tau_2$ ) of V<sub>O</sub>-STFe<sub>0.3</sub> and V<sub>O</sub>-STCo<sub>0.3</sub> were 4.14 ns and 5.43 ns, respectively, which are lower than that of V<sub>O</sub>-STMn<sub>0.2</sub>. This indicates that V<sub>O</sub>-STMn<sub>0.2</sub> has superior optimization effect on the photogenerated carrier separation, since plentiful surface oxygen vacancies act as traps to facilitate the separation of photogenerated carriers [45,46]. These properties lead to a significant improvement in the photocatalytic performance of V<sub>O</sub>-STMn<sub>0.2</sub>. It is worth noting that a speak and tail can be observed on the photocurrent curve (Fig. 5g and Fig. S16b). Qualitatively, this phenomenon can be understood as follows. When the light source is turned on, the semiconductor absorbs photons with energies equal to or greater than the bandgap energy, leading to the rapid generation of electron-hole pairs on a picosecond timescale, resulting in an instantaneous increase in photocurrent [47–49]. Over time, the trapping and re-emission processes of photo-generated charge carriers on the traps gradually come into effect, which causes the current to gradually decay to the equilibrium current [50,51]. When the light is turned off, the slow decay of the current can be understood as a result of the gradual re-emission and recombination of charge carriers from the traps, while the negative current is attributed to electrons continuing to move towards the surface to recombine with the remaining holes [52,53].

### 3.4. Detection of reaction intermediates and CO<sub>2</sub> reduction mechanism

To reveal the mechanism of photocatalytic CO<sub>2</sub> reduction, in situ DRIFTS experiments were carried out on STMn<sub>0.2</sub> and V<sub>O</sub>-STMn<sub>0.2</sub> under illumination. As for V<sub>O</sub>-STMn<sub>0.2</sub> (Fig. 6a), the new infrared absorption peaks at 1615 cm<sup>-1</sup> and 1278 cm<sup>-1</sup> originate from the C=O and O-H stretching of COOH\*, respectively, which are regarded as critical

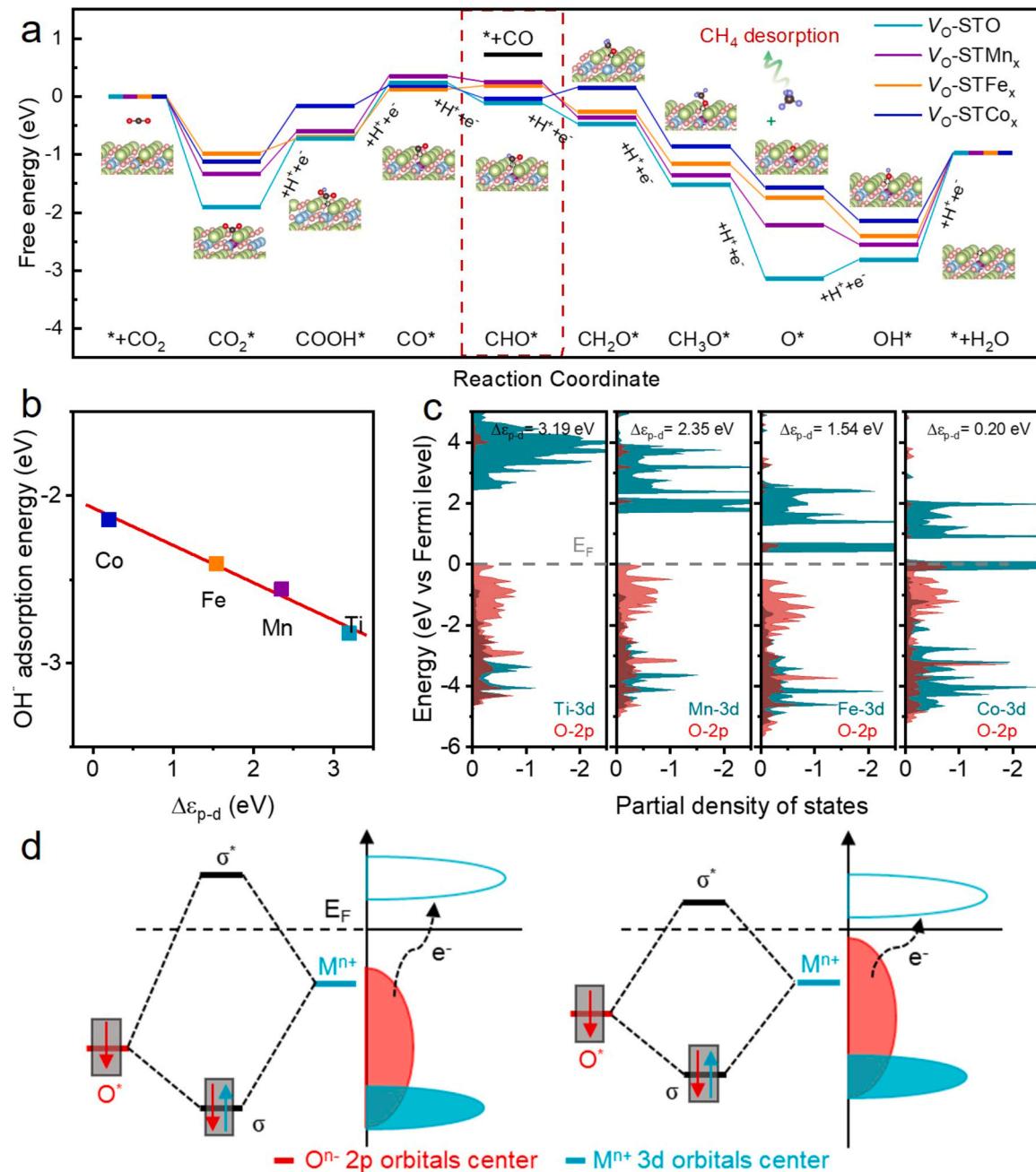


**Fig. 6.** In situ DRIFTS test for  $\text{CO}_2$  and  $\text{H}_2\text{O}$  interaction with (a)  $\text{V}_0\text{-STMn}_{0.2}$  and (b)  $\text{STMn}_{0.2}$  under constant Xenon lamp illumination. Amplifications of the areas from  $2070\text{ cm}^{-1}$  to  $2090\text{ cm}^{-1}$  in (a) and (b), in which the peak at  $2077\text{ cm}^{-1}$  corresponds to the adsorbed  $\text{CO}^*$  molecule.

intermediates during the  $\text{CO}_2$  conversion to  $\text{CO}$  or  $\text{CH}_4$  [54,55]. The intensity of these peaks then progressively increased with the extension of the illumination time. Further, the peak located at  $1110\text{ cm}^{-1}$  is associated with the  $\text{H}-\text{C}=\text{O}$  bending vibration of the  $\text{CHO}^*$  species, and the peaks at  $1135\text{ cm}^{-1}$  and  $1156\text{ cm}^{-1}$  can be attributed to the  $\text{CH}_3\text{O}^*$  species [56]. Both  $\text{CHO}^*$  species and  $\text{CH}_3\text{O}^*$  species are important intermediates in the formation of  $\text{CH}_4$ . Interestingly, the intensity of the observed intermediate species on  $\text{V}_0\text{-STMn}_{0.2}$  was significantly higher than that on  $\text{STMn}_{0.2}$  (Fig. 6b), further demonstrating the excellent  $\text{CO}_2$  methanation activity of  $\text{V}_0\text{-STMn}_{0.2}$ . Another drastic distinction between  $\text{STMn}_{0.2}$  and  $\text{V}_0\text{-STMn}_{0.2}$  arises in the  $\text{CO}^*$  species absorption band at  $\sim 2077\text{ cm}^{-1}$ . From the enlarged region ( $2070\text{--}2090\text{ cm}^{-1}$ ) of Figs. 6a and 6b, the peak intensity of  $\text{CO}^*$  for  $\text{STMn}_{0.2}$  increased over the whole time, while for  $\text{V}_0\text{-STMn}_{0.2}$  the peak intensity of  $\text{CO}^*$  showed an initial increase and then gradually decreased [57]. It is indicated that  $\text{CO}^*$  species on the  $\text{V}_0\text{-STMn}_{0.2}$  surface can be rapidly protonated to  $\text{CHO}^*$  and  $\text{CH}_3\text{O}^*$  intermediates instead of being desorbed to form  $\text{CO}$  molecules, resulting in a nearly 100 % product selectivity for the conversion of  $\text{CO}_2$  to  $\text{CH}_4$ . In this way, the photocatalytic  $\text{CO}_2$  reduction processes for different samples could be described as follows:



The Gibbs free energy of the proposed reduction pathway was calculated to investigate the effect of low-coordination transition metal sites on the  $\text{CO}_2$  reduction activity (Fig. 7a, Figs. S17 and S18). The overall reaction energy barrier for all oxygen vacancy-containing perovskites is lower than that for oxygen vacancy-free perovskites. The introduction of oxygen vacancies leads to charge enrichment on the low coordination transition metal sites, which allows these metal sites to effectively adsorb and activate the reactants. The desorption of  $\text{CO}^*$  on  $\text{STM}_x$  is more favorable than the hydrogenation step, thus  $\text{CO}$  is the major product for  $\text{STM}_x$ . Comparatively,  $\text{V}_0\text{-STM}_x$  preferentially hydrogenates  $\text{CO}^*$ , enabling the selectivity for the  $\text{CH}_4$  product close to 100 % (Fig. 4a). Importantly, the rate-determining step (RDS) for all oxygen vacancy-containing perovskites is the desorption of  $\text{OH}^*$  group. The free energy of oxygen-intermediates ( $\text{CH}_2\text{O}^*$ ,  $\text{CH}_3\text{O}^*$ ,  $\text{O}^*$ ,  $\text{OH}^*$ ) with the O atom as the center of interaction for  $\text{V}_0\text{-STO}$  is the most negative of all  $\text{V}_0\text{-STM}_x$  perovskites studied, indicating that the Ti active sites of  $\text{V}_0\text{-STO}$  have the strongest adsorption for these oxygen-intermediates. The RDS barrier of  $\text{V}_0\text{-STM}_x$  is lower than that of  $\text{V}_0\text{-STO}$ , and the limiting barriers of the perovskites  $\text{V}_0\text{-STO}$ ,  $\text{V}_0\text{-STM}_x$ ,  $\text{V}_0\text{-STFex}$  and



**Fig. 7.** (a) The Gibbs free energy diagram of  $\text{CO}_2$  photoreduction to  $\text{CH}_4$  for the perovskites  $\text{V}_0\text{-STO}$  and  $\text{V}_0\text{-STM}_x$ . (b) The linear relationship for the adsorption energy of  $\text{OH}^*$  group and the energy difference ( $\Delta\epsilon_{p-d}$ ) between the O atom 2p band center and the transition metal cation d band center. (c) The partial density of states for the 3d orbitals of transition metal atom and the O 2p orbitals of  $\text{OH}^*$  group. The band centers of spin down were used. (d) Schematic illustration for the hybridization of the d orbital of the metal cation with the oxygen p orbital of the adsorbate with the decrement of  $\Delta\epsilon_{p-d}$ .

$\text{V}_0\text{-STCo}_x$  are 1.84, 1.58, 1.47 and 1.17 eV, respectively, suggesting that the doped transition metal ions weaken the interaction with the oxygen-intermediates.

According to the crystal field effect, the static electric field splits the d orbitals of the  $\text{MO}_6$  octahedral metal cation into three low-lying  $t_{2g}$  orbitals ( $d_{xy}$ ,  $d_{xz}$ ,  $d_{yz}$ ) and two high-lying  $e_g$  orbitals ( $d_{x^2-z^2}$ ,  $d_{z^2}$ ) [58]. The split d orbitals of the metal cation contribute to the narrow d band. The principle behind the linear relationship for the adsorption energy of the oxygen-intermediates and the energy difference ( $\Delta\epsilon_{p-d}$ ) between the O atom 2p band center and the metal cation d band center is due to the coupling of the metal cation d orbitals and the O atom 2p orbitals (Figs. 7b and 7c). When low-coordination transition metal cation sites interact with oxygen intermediates, the electrons of transition metal

cation d orbitals and 2p orbitals of O atoms spatially overlap, forming  $\sigma$  bonding molecular orbital states and  $\sigma^*$  antibonding molecular orbital states (Fig. 7d). As the 2p orbitals center of the O atom approaches or rises above the d orbitals center of the metal cation (smaller  $\Delta\epsilon_{p-d}$ ), the unbonded electrons of the O atom can transition to the antibonding orbital with relatively less energy. The reduced charge of the O atom and the decreased energy of the antibonding orbitals give the oxygen intermediate a higher degree of freedom [59]. Thus, the Mn atom with a relatively small  $\Delta\epsilon_{p-d}$  may have a more suitable adsorption of oxygen-intermediates, which reduces the unfavorable strong adsorption of oxygen-intermediates and possesses a higher catalytic activity. Another remarkable factor is the influence of the surface oxygen vacancy contents on the  $\text{CO}_2$  reduction properties. In fact, when we take

the oxygen vacancy content as a variable, the  $\text{CO}_2$  reduction performance is enhanced with the increasing oxygen vacancy concentration and follows a good linear relationship, as illustrated in Fig. S19. However, it should be noted that the catalyst only has the best catalytic activity when the vacancy concentration is at an appropriate level [60].

Briefly, the relative center between the O 2p orbitals and the transition metal 3d orbitals could regulate the adsorption ability of the low-coordination metal cation sites for oxygen-intermediates. Making full use of the parameter  $\Delta\epsilon_{\text{p-d}}$ , the mechanism of  $\text{CO}_2$  reduction reaction for a given  $\text{STM}_x$  perovskite can be effectively predicted. Meanwhile, the  $\text{CO}_2$  reduction performance of strontium titanate-based perovskites also depends on the adsorption properties of other oxygen intermediates ( $\text{COOH}^*$ ,  $\text{CO}^*$ ,  $\text{CHO}^*$ ) during  $\text{CO}_2$  reduction, which requires full experimental and theoretical verification in the future.

#### 4. Conclusion

In summary, a model of low-coordination metal active sites was constructed for the tunable selectivity and reactivity of  $\text{CO}_2$  photochemical reduction based on strontium titanate perovskites. After activity modification of the  $\text{STM}_x$  ( $x = 0.1\text{--}0.3$ ) perovskites by reduction with  $\text{NaBH}_4$ , the  $\text{V}_0\text{-STMn}_{0.2}$  exhibited a gas yield of  $18.21 \mu\text{mol g}^{-1}$  with near 100 % selectivity for  $\text{CH}_4$ . The electronic structure analysis suggest that the introduction of oxygen vacancies leads to charge enrichment on the low-coordination Mn atom. In situ DRIFTS and Gibbs free energy calculations reveal that the charge-rich low-coordination metal sites have lowered the overall activation energy barrier, thus facilitating the hydrogenation of  $\text{CO}^*$  to  $\text{CH}_4$ . Theoretical calculations demonstrate that the relative centers between the O atom 2p orbitals and the metal 3d orbitals regulate the adsorption properties of the low-coordination metal sites for oxygen-intermediates. In future work, we will carry out a comprehensive study on A-site modulation of perovskites to explore the design options for bifunctional sites photocatalysts with enhanced adsorption properties for  $\text{COOH}^*$ ,  $\text{CO}^*$  and  $\text{CHO}^*$  to improve the photocatalytic performance. The results of the study provide insight into the design and defective modulation of photocatalytic materials.

#### CRediT authorship contribution statement

**Yibo Gao:** Methodology, Data curation, Formal analysis, Writing – original draft. **Miaomiao Zhang:** Writing – review & editing. **Yang Jin:** Funding acquisition, Writing - review & editing. **Meng Zhou:** Software. **Yanpeng Mao:** Funding acquisition, Conceptualization, Investigation, Resources. **Jian Sun:** Writing – review & editing. **Wenlong Wang:** Resources. **Zhanlong Song:** Resources.

#### Declaration of Competing Interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

#### Data availability

Data will be made available on request.

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#### Appendix A. Supporting information

Supplementary data associated with this article can be found in the online version at [doi:10.1016/j.apcatb.2023.123348](https://doi.org/10.1016/j.apcatb.2023.123348).

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